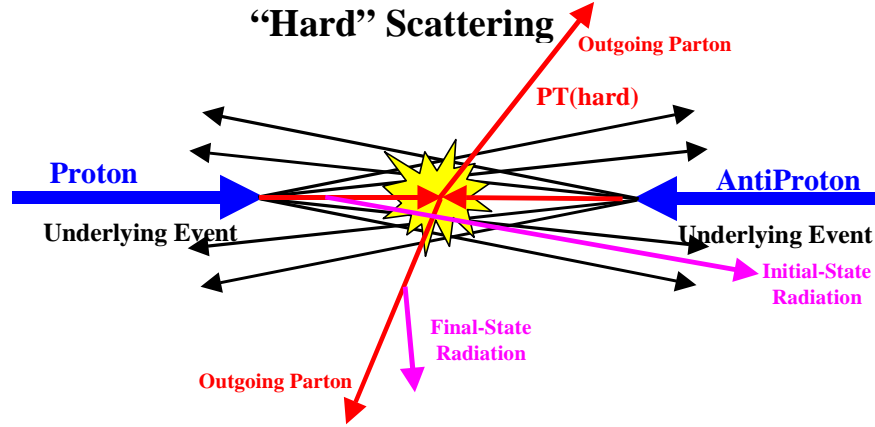
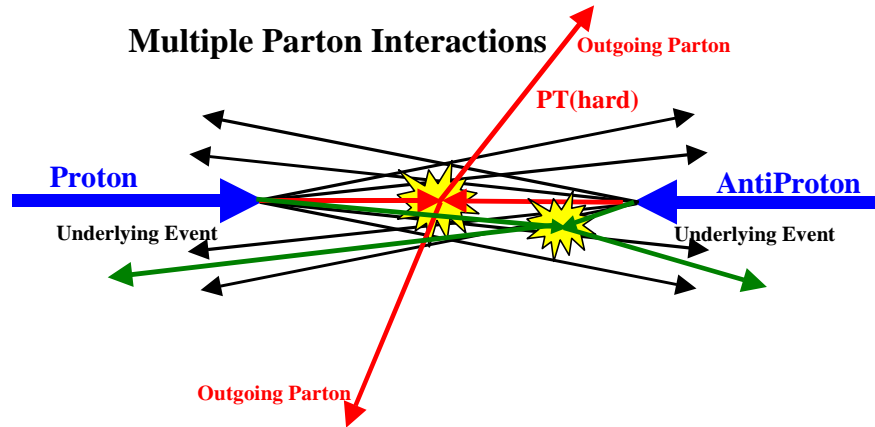


# The Underlying Event: Z-boson vs Dijet

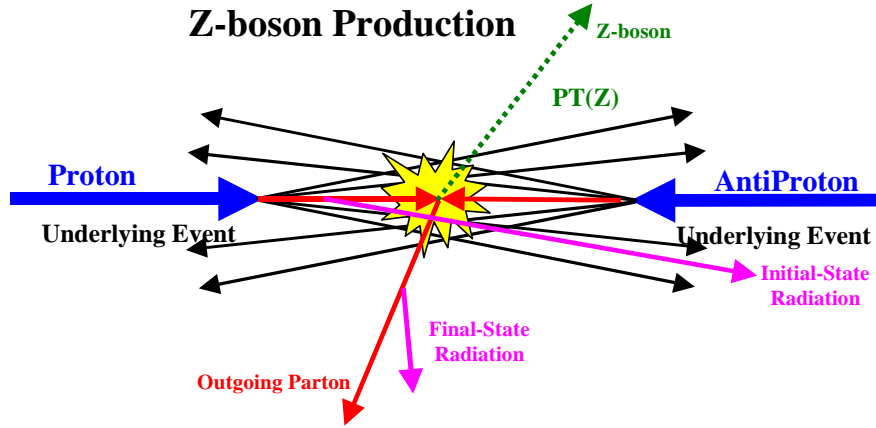
Rick Field, Henry Frisch, Richard Haas, and David Stuart



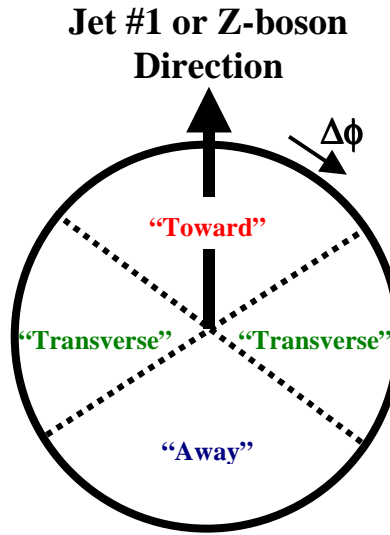
**Fig. 1.** Illustration of a proton-antiproton collision in which a “hard” 2-to-2 parton scattering with transverse momentum,  $P_T(\text{hard})$ , has occurred. The resulting “dijet” event contains particles that originate from the two outgoing partons (plus final-state radiation) and particles that come from the breakup of the proton and antiproton (*i.e.* “beam-beam remnants”). The “underlying event” consists of the beam-beam remnants plus initial-state radiation.



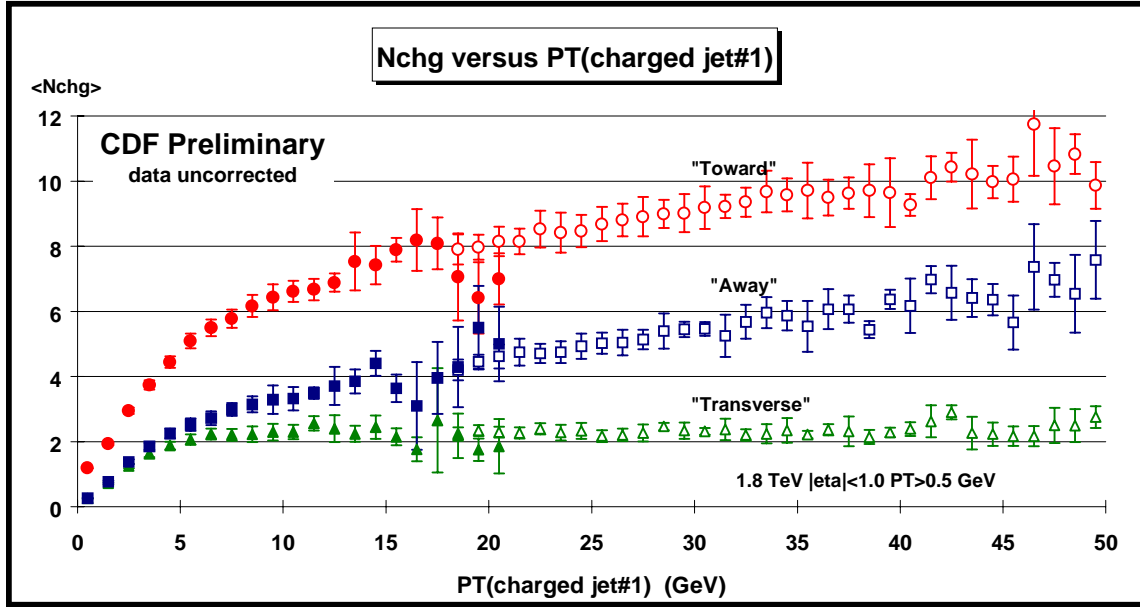
**Fig. 2.** Illustration of a proton-antiproton collision in which a multiple parton interaction has occurred. In addition to the “hard” 2-to-2 parton scattering with transverse momentum,  $P_T(\text{hard})$ , there is an additional “semi-hard” parton-parton scattering that contributes particles to the “underlying event”. For PYTHIA, we include the contributions from multiple parton scattering in the beam-beam remnant component.



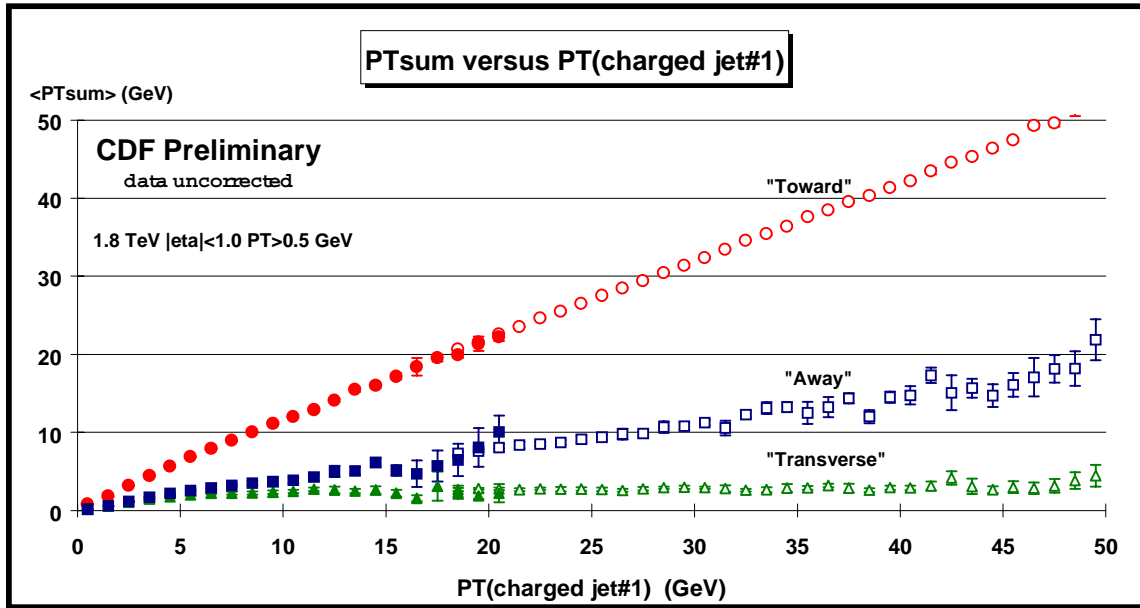
**Fig. 3.** Illustration of a proton-antiproton collision in which a Z-boson with large transverse momentum,  $P_T(Z)$ , has been produced. The resulting event contains particles that originate from the “away-side” outgoing parton (plus final-state radiation) and particles that come from the breakup of the proton and antiproton (*i.e.* “beam-beam remnants”). The “underlying event” consists of the beam-beam remnants plus initial-state radiation.



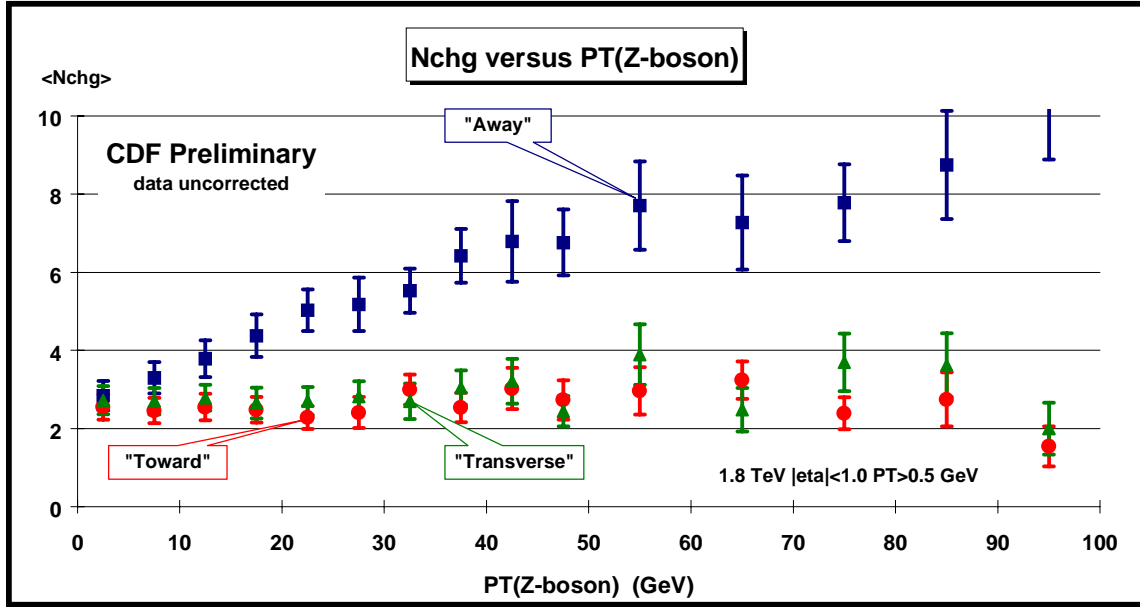
**Fig. 4.** Illustration of correlations in azimuthal angle  $\Delta\phi$  relative to the direction of the leading charged jet in the event, jet#1, or the Z-boson. The angle  $\Delta\phi = |\phi - \phi_{\text{jet}\#1}|$  is the relative azimuthal angle between charged particles and the direction of jet#1 or the Z-boson. The region  $|\Delta\phi| < 60^\circ$  is referred to as “toward” (*includes particles in jet#1 but does not include the decay products of the Z-boson*) and the region  $|\Delta\phi| > 120^\circ$  is called “away”. The “transverse” region is defined by  $60^\circ < |\Delta\phi| < 120^\circ$ . Each region, “toward”, “transverse”, and “away” covers the same range  $|\Delta\eta| \times |\Delta\phi| = 2 \times 120^\circ$ .



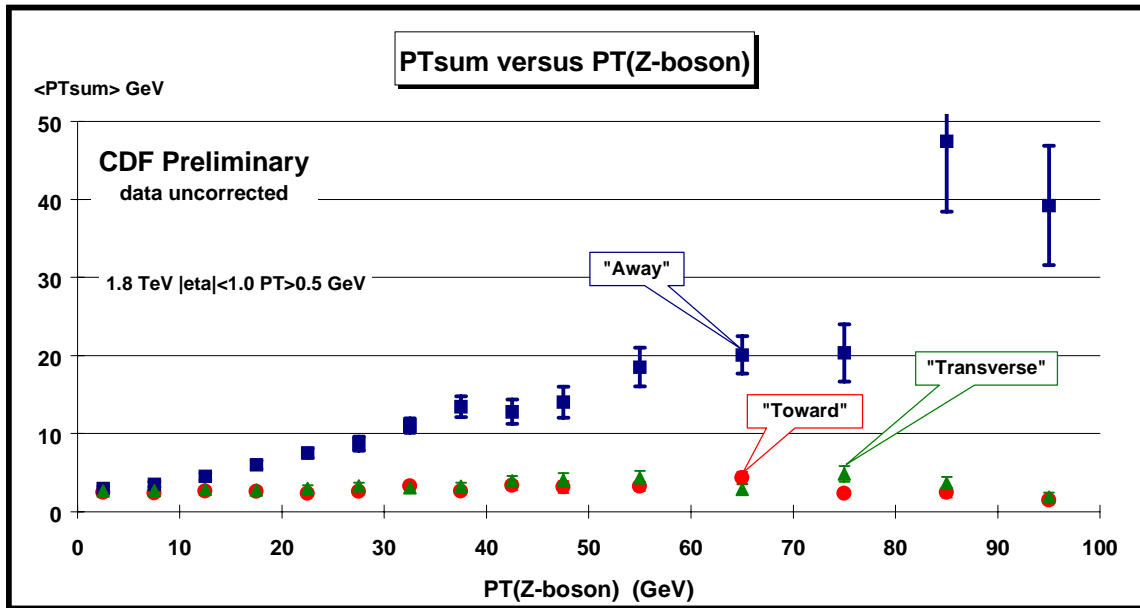
**Fig. 5.** The average number of “toward” ( $|\Delta\phi| < 60^\circ$ ), “transverse” ( $60 < |\Delta\phi| < 120^\circ$ ), and “away” ( $|\Delta\phi| > 120^\circ$ ) charged particles ( $P_T > 0.5$  GeV,  $|\eta| < 1$ , including jet#1) as a function of the transverse momentum of the leading charged particle jet. Each point corresponds to the  $\langle N_{\text{chg}} \rangle$  in a 1 GeV bin. The solid (open) points are the Min-Bias (JET20) data. The errors on the (uncorrected) data include both statistical and correlated systematic uncertainties. The “toward”, “transverse”, and “away” regions are defined in Fig. 4.



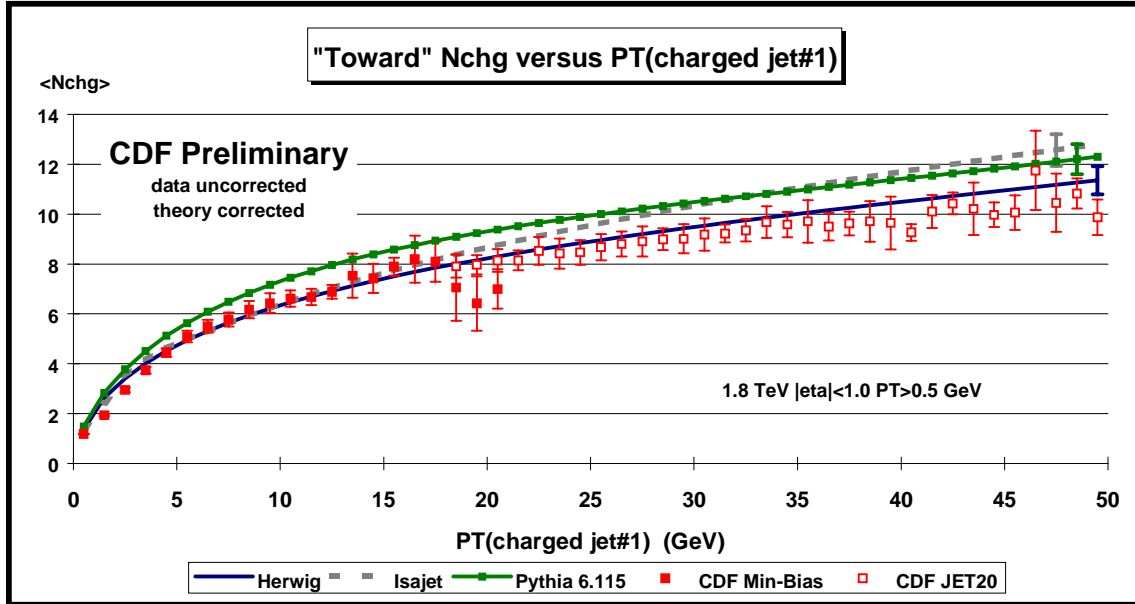
**Fig. 6.** The average scalar  $P_T$  sum of “toward” ( $|\Delta\phi| < 60^\circ$ ), “transverse” ( $60 < |\Delta\phi| < 120^\circ$ ), and “away” ( $|\Delta\phi| > 120^\circ$ ) charged particles ( $P_T > 0.5$  GeV,  $|\eta| < 1$ , including jet#1) as a function of the transverse momentum of the leading charged particle jet. Each point corresponds to the  $\langle P_{T\text{sum}} \rangle$  in a 1 GeV bin. The solid (open) points are the Min-Bias (JET20) data. The errors on the (uncorrected) data include both statistical and correlated systematic uncertainties. The “toward”, “transverse”, and “away” regions are defined in Fig. 4.



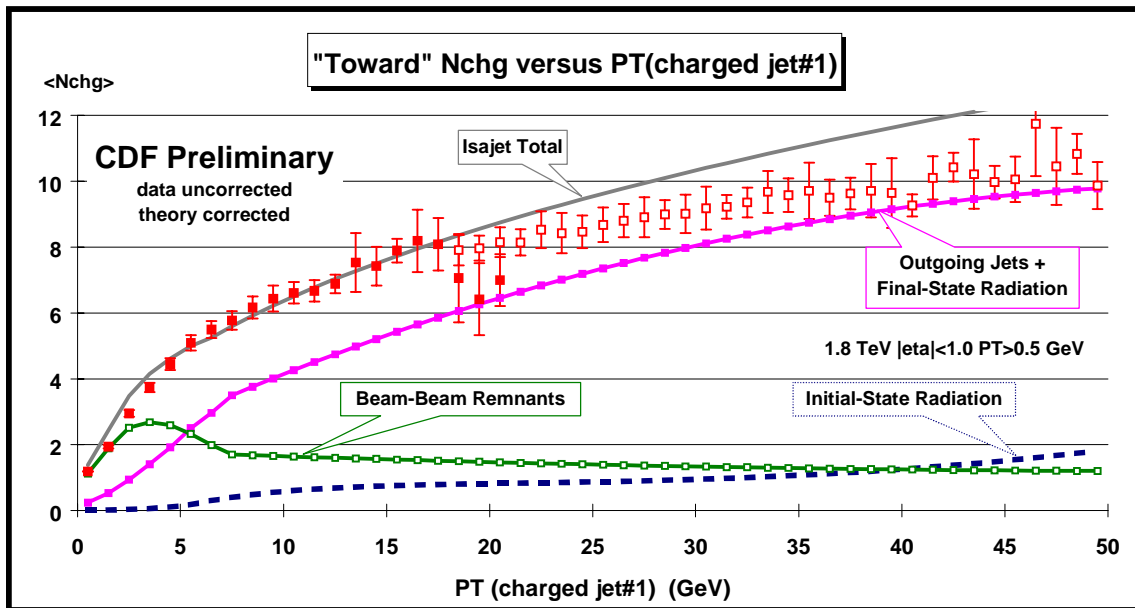
**Fig. 7.** The average number of “toward” ( $|\Delta\phi| < 60^\circ$ ), “transverse” ( $60 < |\Delta\phi| < 120^\circ$ ), and “away” ( $|\Delta\phi| > 120^\circ$ ) charged particles ( $P_T > 0.5$  GeV,  $|\eta| < 1$ , excluding decay products of the Z-boson) as a function of the transverse momentum of the Z-boson. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The “toward”, “transverse”, and “away” regions are defined in Fig. 4.



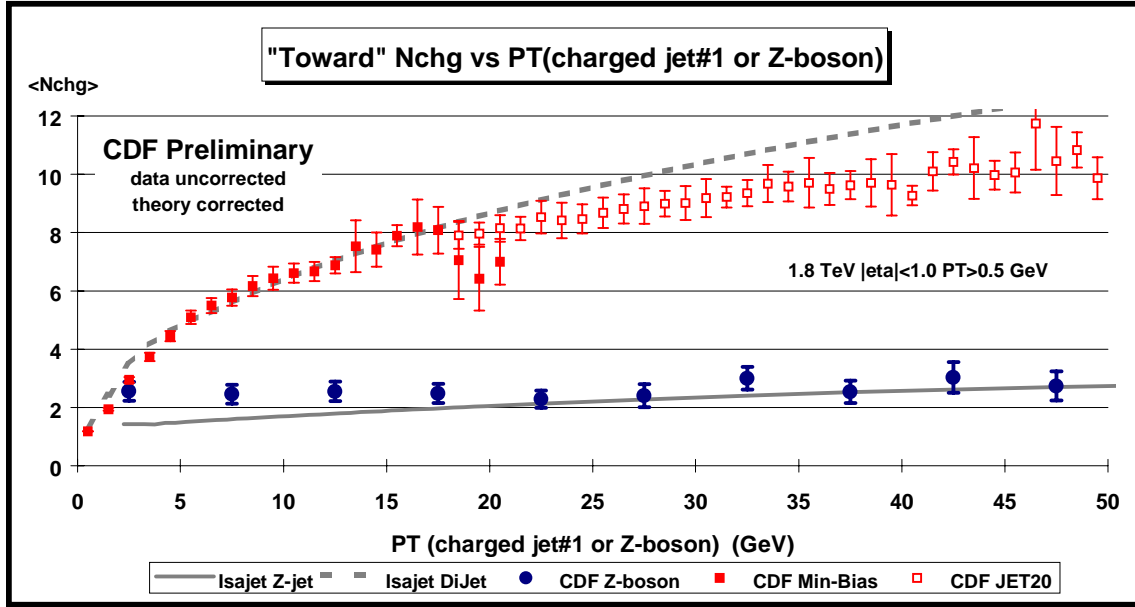
**Fig. 8.** The average *scalar*  $P_T$  sum of “toward” ( $|\Delta\phi| < 60^\circ$ ), “transverse” ( $60 < |\Delta\phi| < 120^\circ$ ), and “away” ( $|\Delta\phi| > 120^\circ$ ) charged particles ( $P_T > 0.5$  GeV,  $|\eta| < 1$ , excluding the decay products of the Z-boson) as a function of the transverse momentum of the Z-boson. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The “toward”, “transverse”, and “away” regions are defined in Fig. 4.



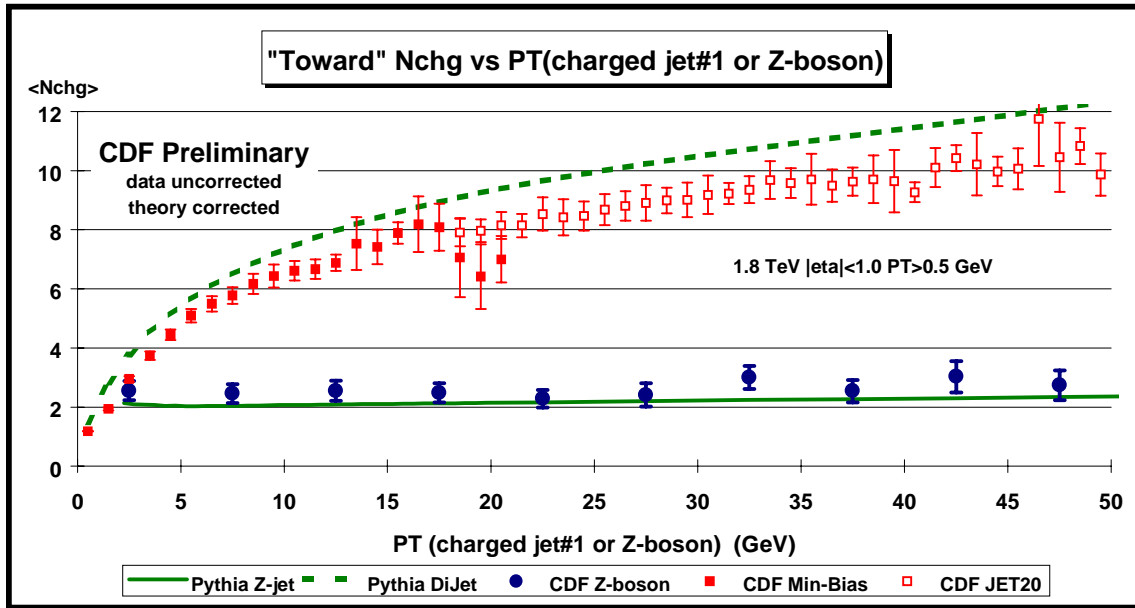
**Fig. 9.** Dijet data from Fig. 5 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet\#1})$  (*leading charged jet*) for the “toward” region defined in Fig. X compared with the QCD “hard scattering” Monte-Carlo predictions of HERWIG 5.9, ISAJET 7.32, and PYTHIA 6.115. Each point corresponds to the “toward”  $\langle N_{\text{chg}} \rangle$  in a 1 GeV bin. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



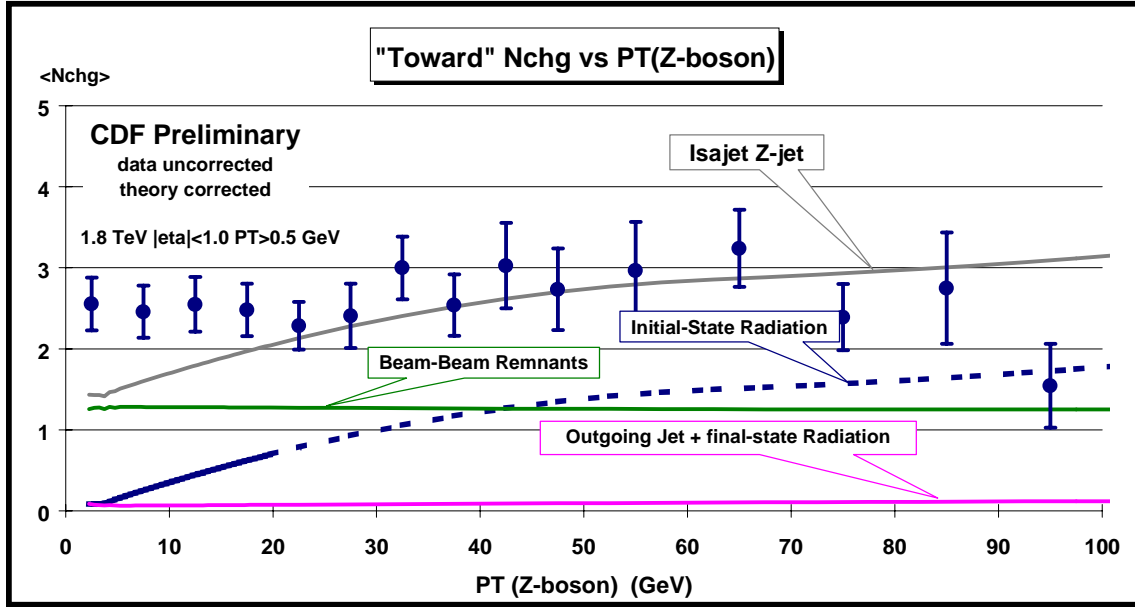
**Fig. 10.** Dijet data from Fig. 5 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet\#1})$  (*leading charged jet*) for the “toward” region defined in Fig. X compared with the QCD “hard scattering” Monte-Carlo predictions of ISAJET 7.32. The predictions of ISAJET are divided into three categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), charged particles that arise from initial-state radiation, and charged particles that result from the outgoing jets plus final-state radiation (see Fig. 1). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



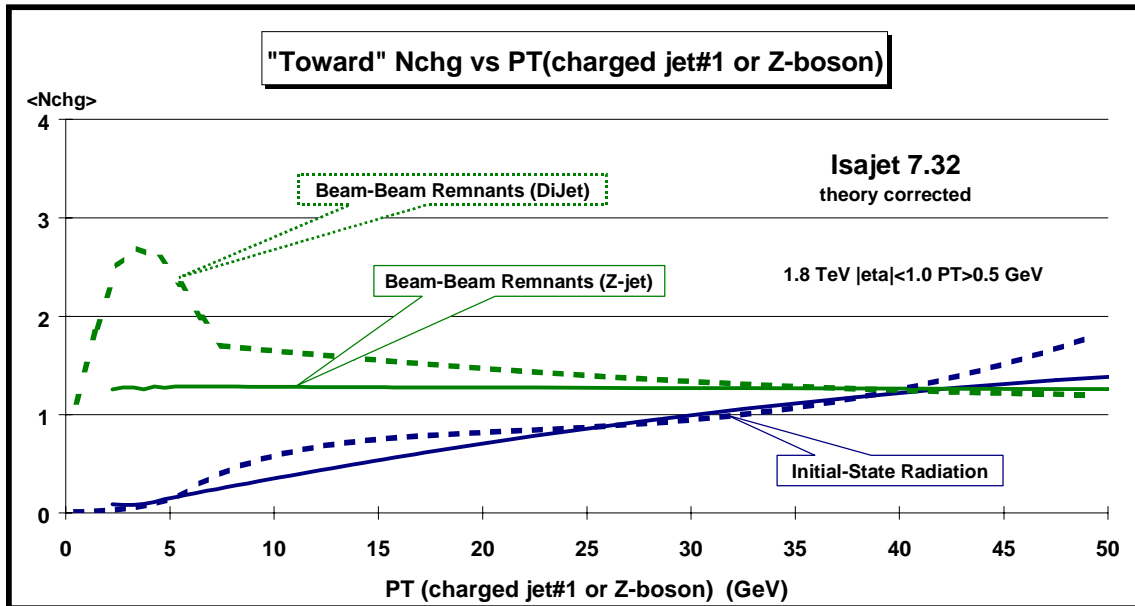
**Fig. 11.** Comparison of the dijet data from Fig. 5 and the Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “toward” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of ISAJET 7.32 for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



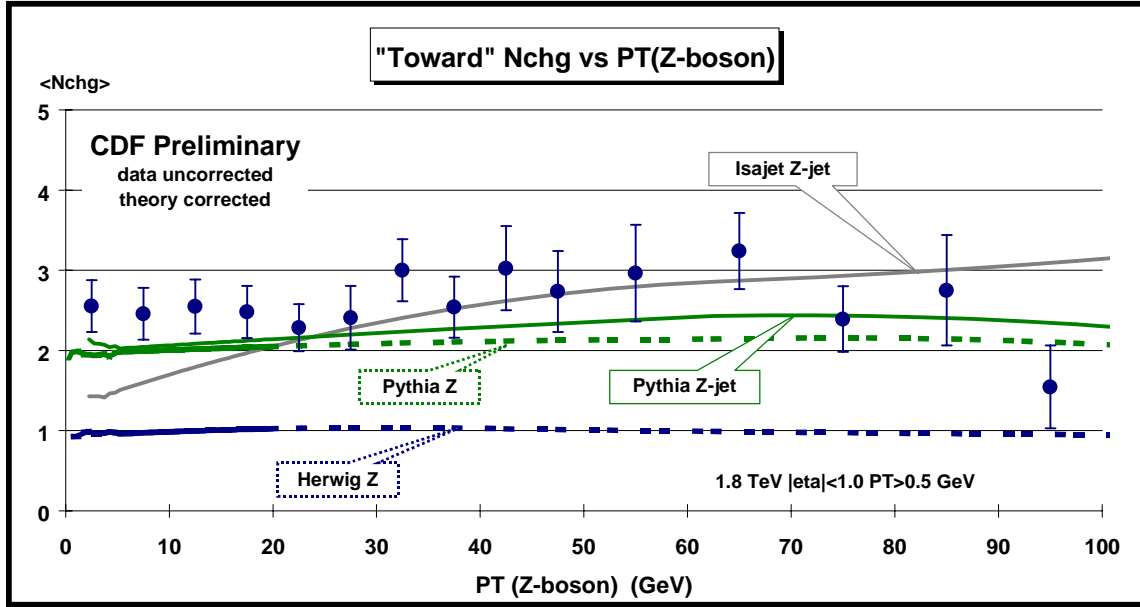
**Fig. 12.** Comparison of the dijet data from Fig. 5 and the Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “toward” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of PYTHIA 6.115 for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



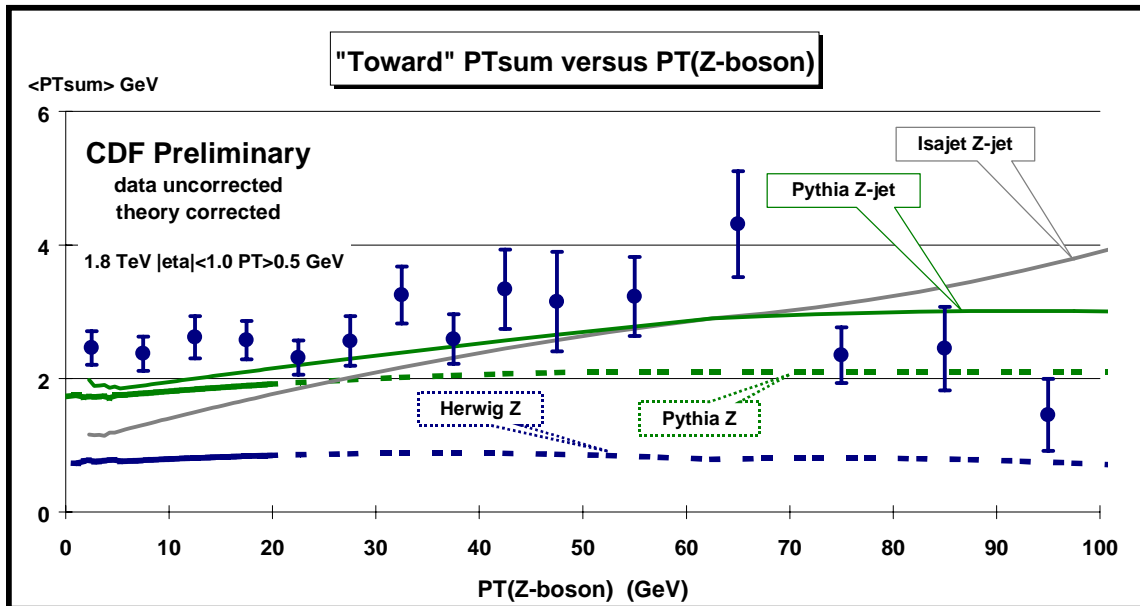
**Fig. 13.** Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “toward” region defined in Fig. 4 compared with the “Z-jet” QCD Monte-Carlo predictions of ISAJET 7.32. The predictions of ISAJET are divided into three categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), charged particles that arise from initial-state radiation, and charged particles that result from the outgoing jets plus final-state radiation (see Fig. 3). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 14.** Comparison of the QCD Monte-Carlo predictions of ISAJET 7.32 for the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “toward” region defined in Fig. 4 for dijet (dashed) and “Z-jet” (solid) production. The plot shows the charged particles that arise from the break-up of the beam and target (*beam-beam remnants*) and the charged particles that arise from from initial-state radiation (see Fig. 1 and Fig. 3). The curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

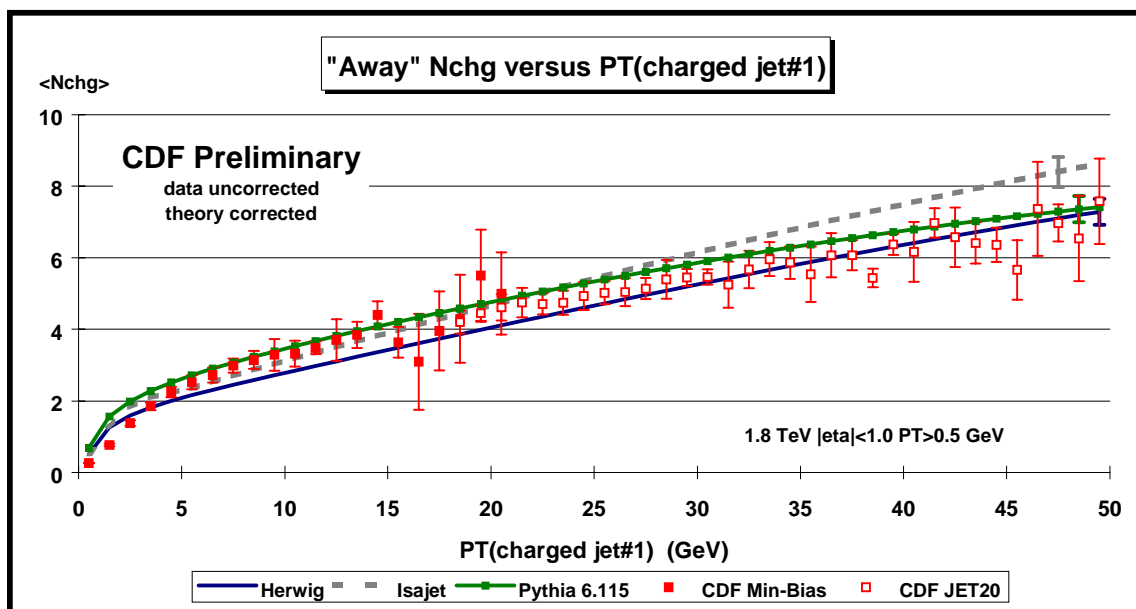


**Fig. 15.** Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “toward” region defined in Fig. 4 compared with the QCD Monte-Carlo predictions of HERWIG 5.9 (“Z”), ISAJET 7.32 (“Z-jet”), and PYTHIA 6.115 (“Z”, “Z-jet”) (see table 3). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

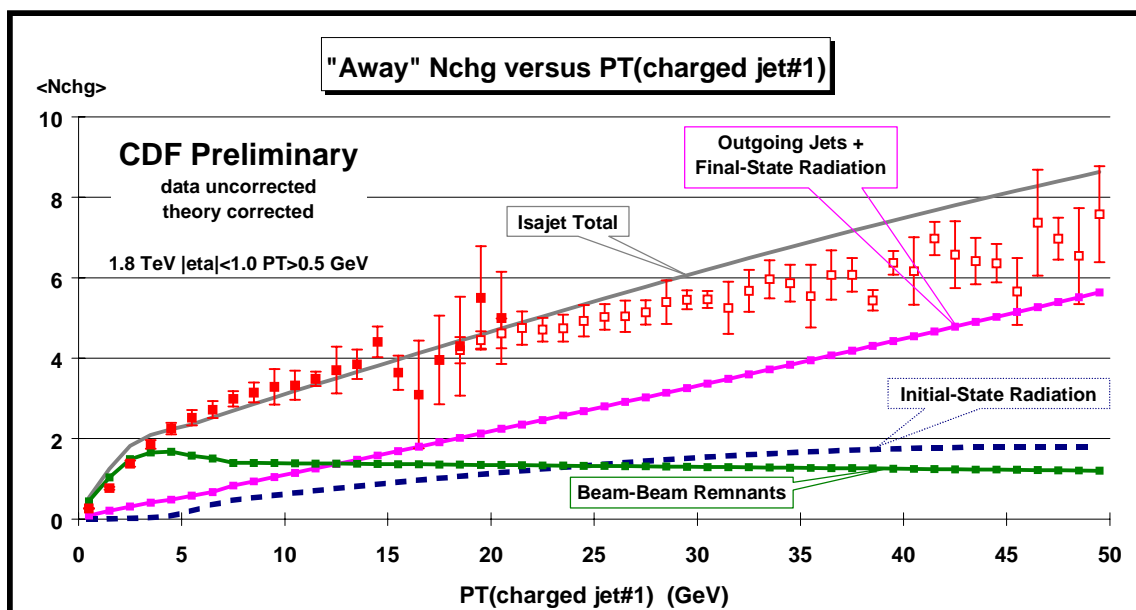


**Fig. 16.** Z-boson data from Fig. 8 on the average *scalar*  $P_T$  sum of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “toward” region defined in Fig. 4 compared with the QCD Monte-Carlo predictions of HERWIG 5.9 (“Z”), ISAJET 7.32 (“Z-jet”), and PYTHIA 6.115 (“Z”, “Z-jet”) (see table 3). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

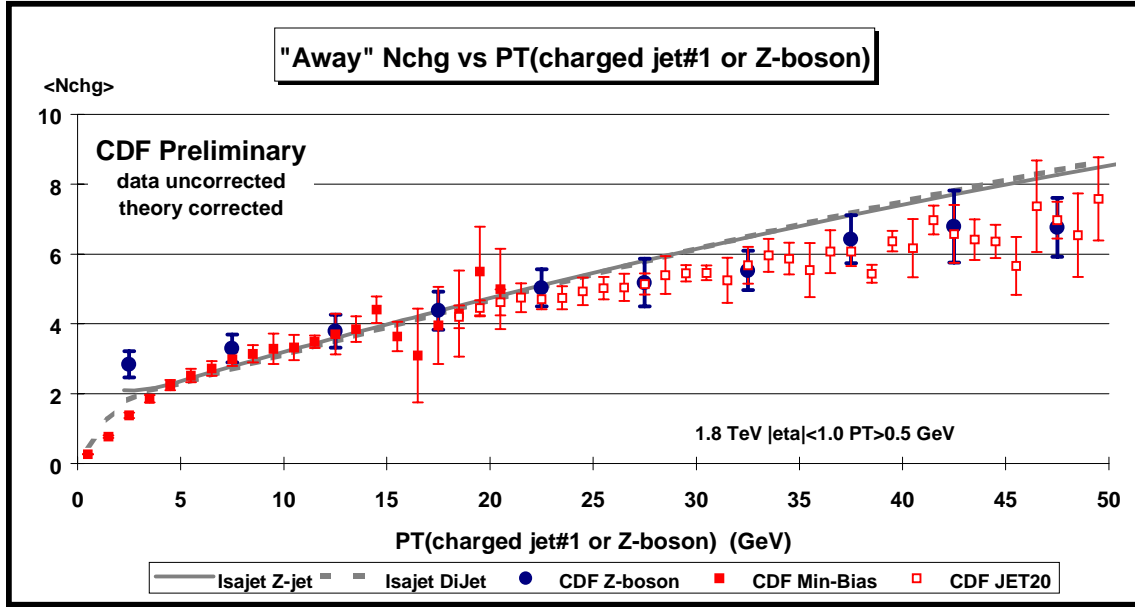




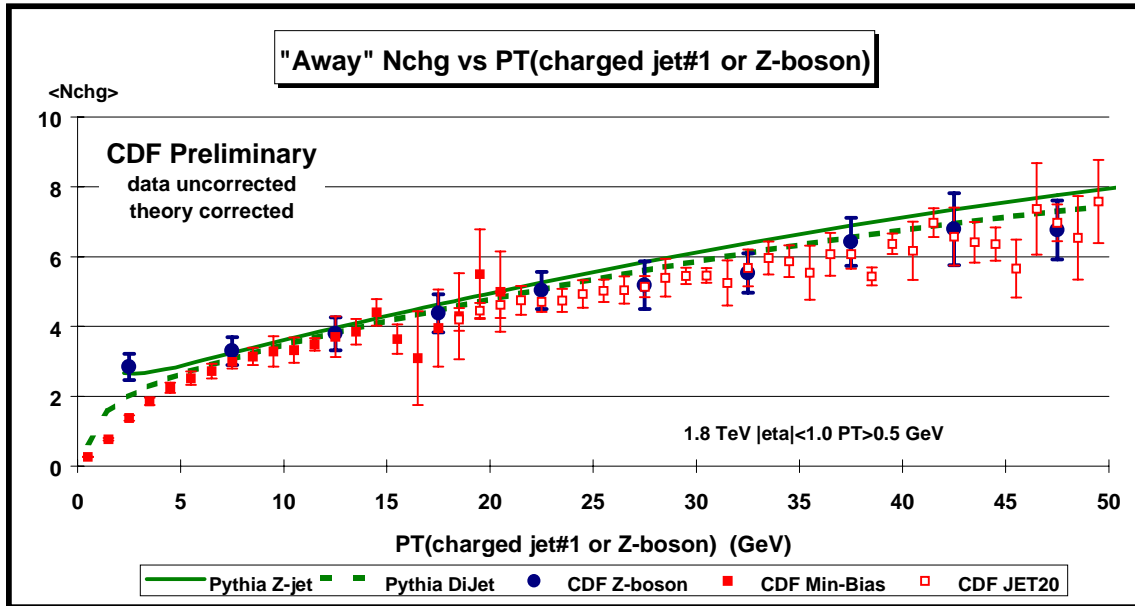
**Fig. 17.** Dijet data from Fig. 5 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet\#1})$  (*leading charged jet*) for the “away” region defined in Fig. 4 compared with the QCD “hard scattering” Monte-Carlo predictions of HERWIG 5.9, ISAJET 7.32, and PYTHIA 6.115. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 18.** Dijet data from Fig. 5 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet\#1})$  (*leading charged jet*) for the “away” region defined in Fig. 4 compared with the QCD “hard scattering” Monte-Carlo predictions of ISAJET 7.32. The predictions of ISAJET are divided into three categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), charged particles that arise from initial-state radiation, and charged particles that result from the outgoing jets plus final-state radiation (see Fig. 1). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 19.** Comparison of the dijet data from Fig. 5 and the Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “away” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of ISAJET 7.32 for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 20.** Comparison of the dijet data from Fig. 5 and the Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “away” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of PYTHIA 6.115 for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

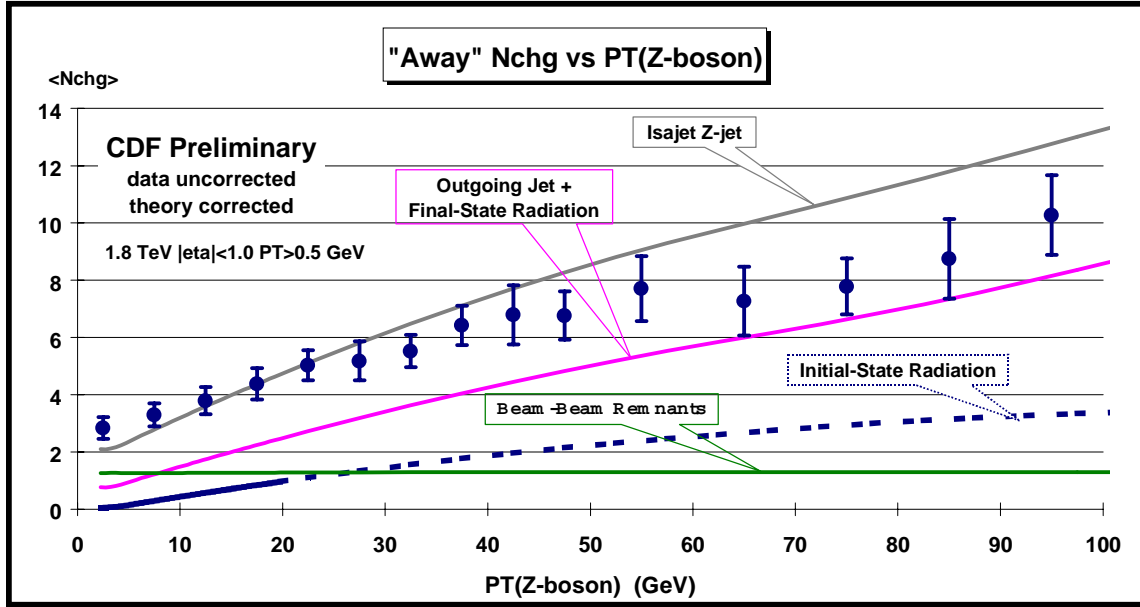


Fig. 21. Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $PT(Z)$  for the “away” region defined in Fig. 4 compared with the “Z-jet” QCD Monte-Carlo predictions of ISAJET 7.32. The predictions of ISAJET are divided into three categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), charged particles that arise from initial-state radiation, and charged particles that result from the outgoing jet plus final-state radiation (see Fig. 3). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

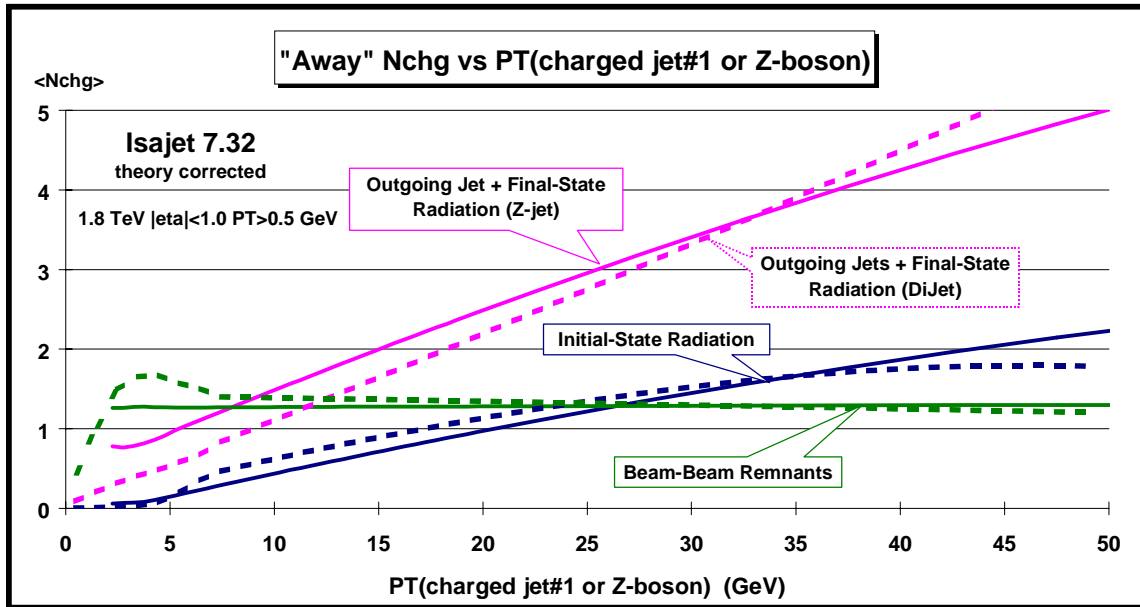
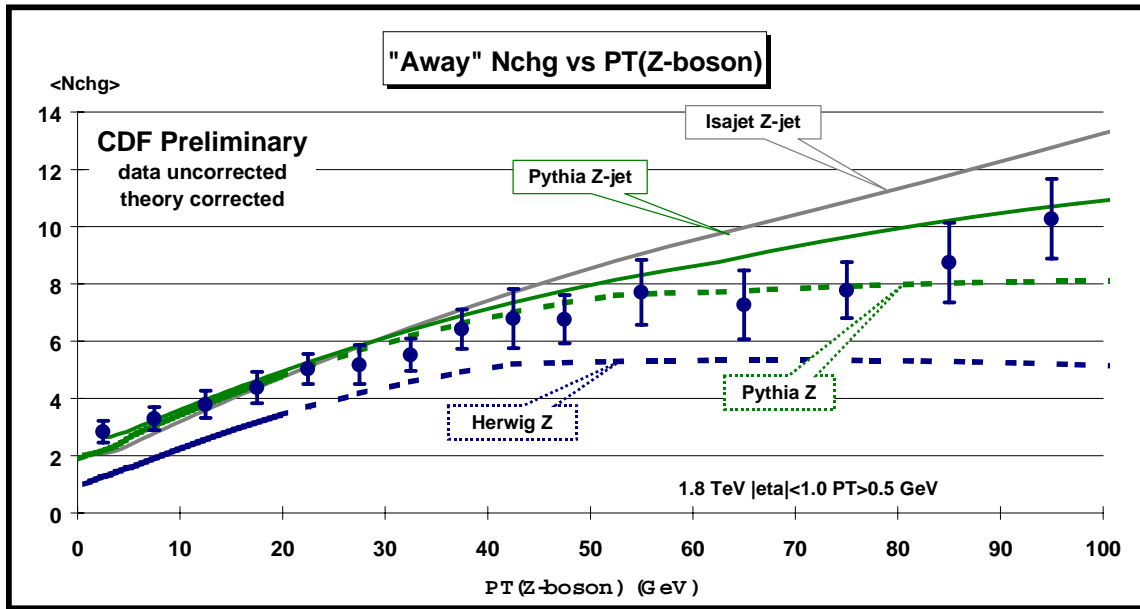
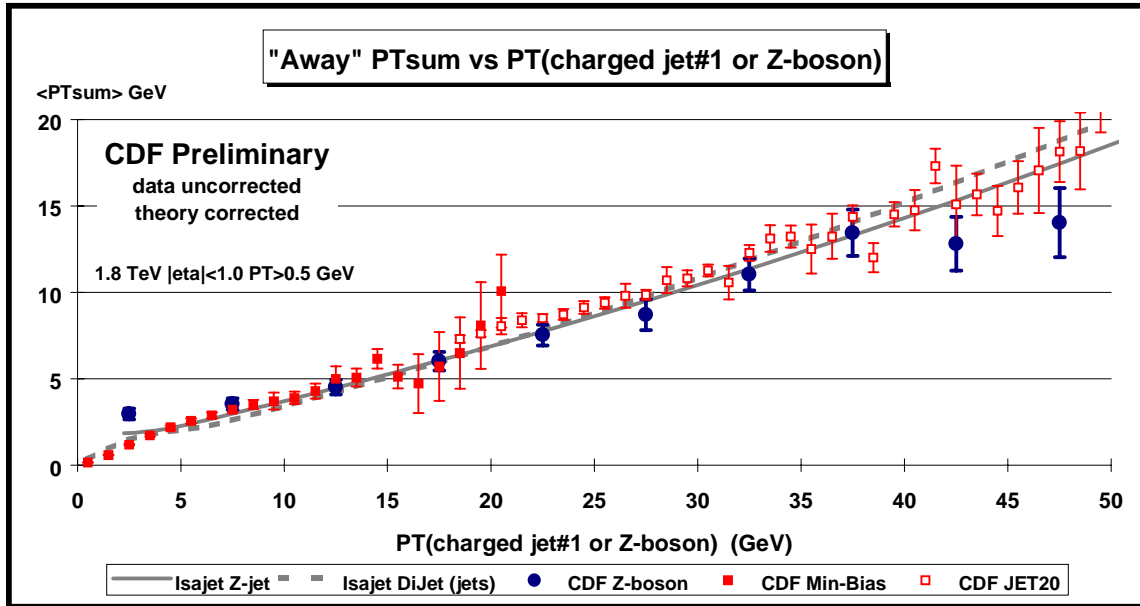


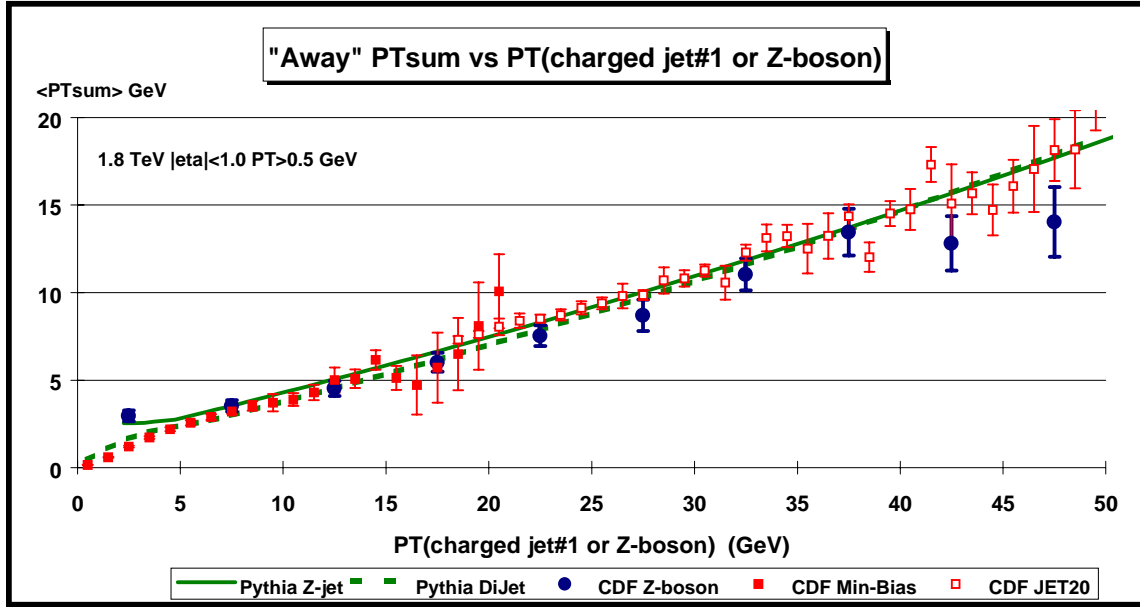
Fig. 22. Comparison of the QCD Monte-Carlo predictions of ISAJET 7.32 for the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “away” region defined in Fig. 4 for dijet Fig. 18 (dashed) and “Z-jet” Fig. 21 (solid) production. The predictions of ISAJET are divided into three categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), charged particles that arise from initial-state radiation, and charged particles that result from the outgoing jet(s) plus final-state radiation (see Fig. 1 and Fig.3). The curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



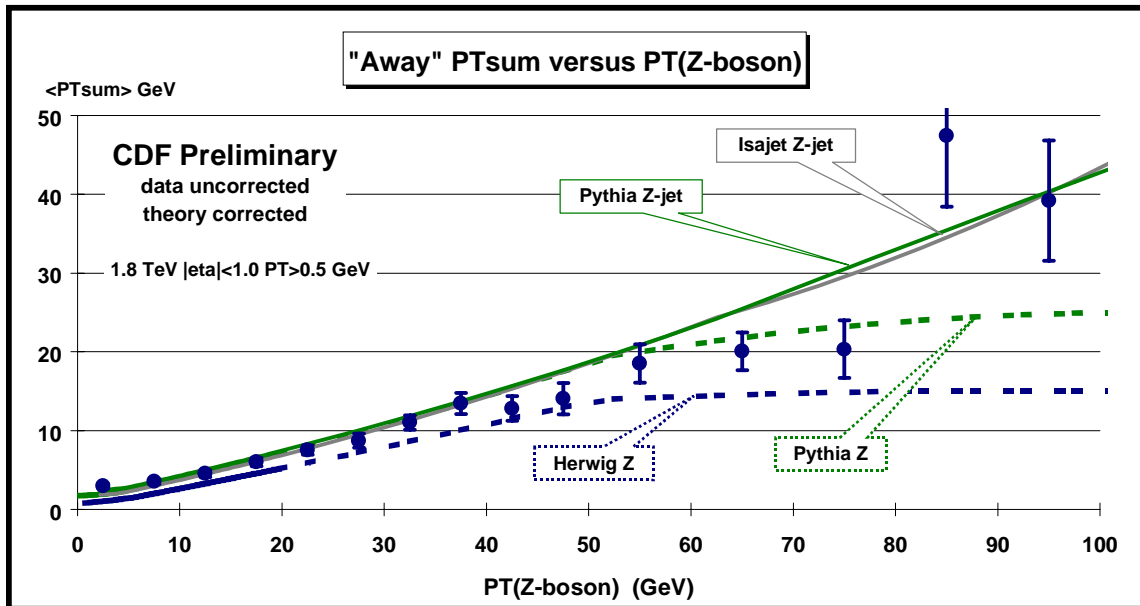
**Fig. 23.** Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “away” region defined in Fig. 4 compared with the QCD Monte-Carlo predictions of HERWIG 5.9 (“Z”), ISAJET 7.32 (“Z-jet”), and PYTHIA 6.115 (“Z”, “Z-jet”) (see table 3). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 24.** Comparison of the dijet data from Fig. 6 and the Z-boson data from Fig. 8 on the average *scalar*  $P_T$  sum of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “away” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of ISAJET 7.32 for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 25.** Comparison of the dijet data from Fig. 6 and the Z-boson data from Fig. 8 on the average *scalar*  $P_T$  sum of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “away” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of PYTHIA 6.115 for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 26.** Z-boson data from Fig. 8 on the average *scalar*  $P_T$  sum of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T$ (Z-boson) for the “away” region defined in Fig. 4 compared with the QCD Monte-Carlo predictions of HERWIG 5.9 (“Z”), ISAJET 7.32 (“Z-jet”), and PYTHIA 6.115 (“Z”, “Z-jet”) (see table 3). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

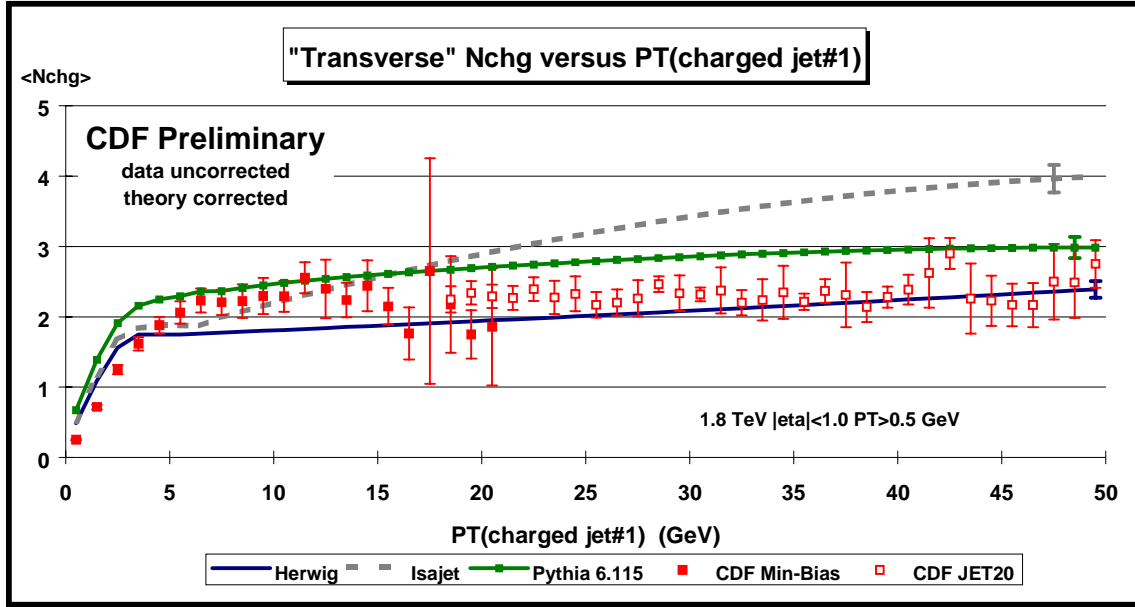


Fig. 27. Dijet data from Fig. 5 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet\#1})$  (leading charged jet) for the “transverse” region defined in Fig. 4 compared with the QCD “hard scattering” Monte-Carlo predictions of HERWIG 5.9, ISAJET 7.32, and PYTHIA 6.115. The errors on the (uncorrected) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (statistical plus systematic) of around 5%.

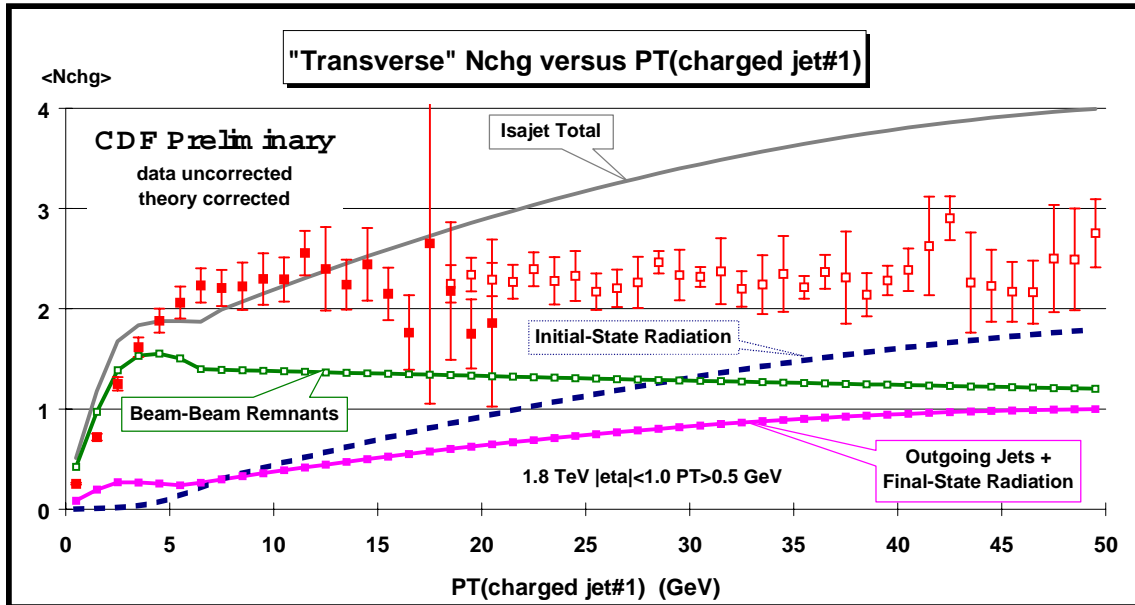
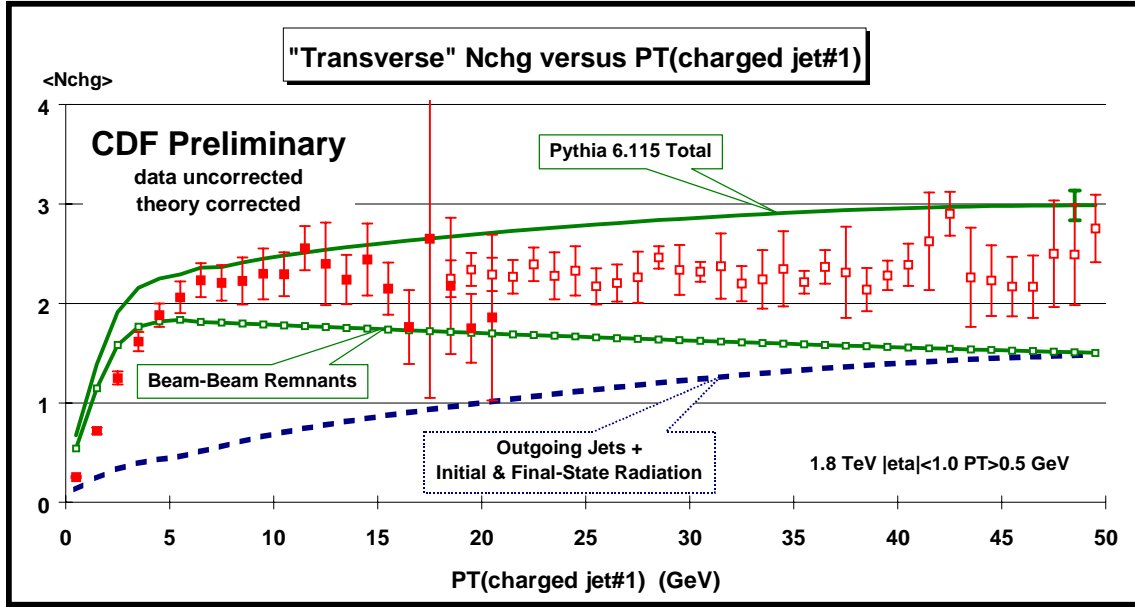
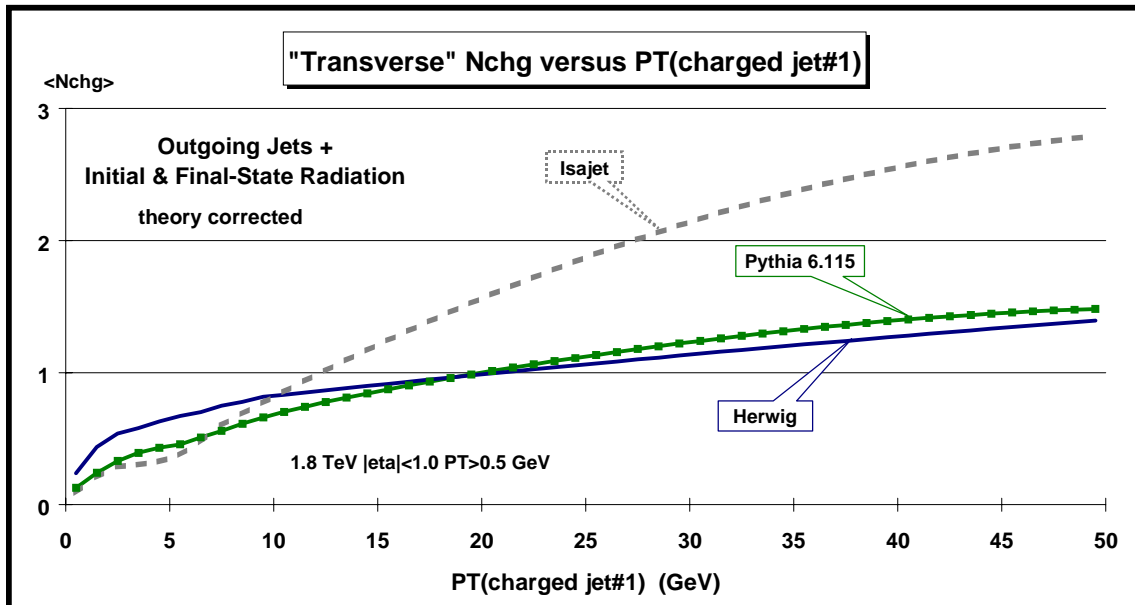


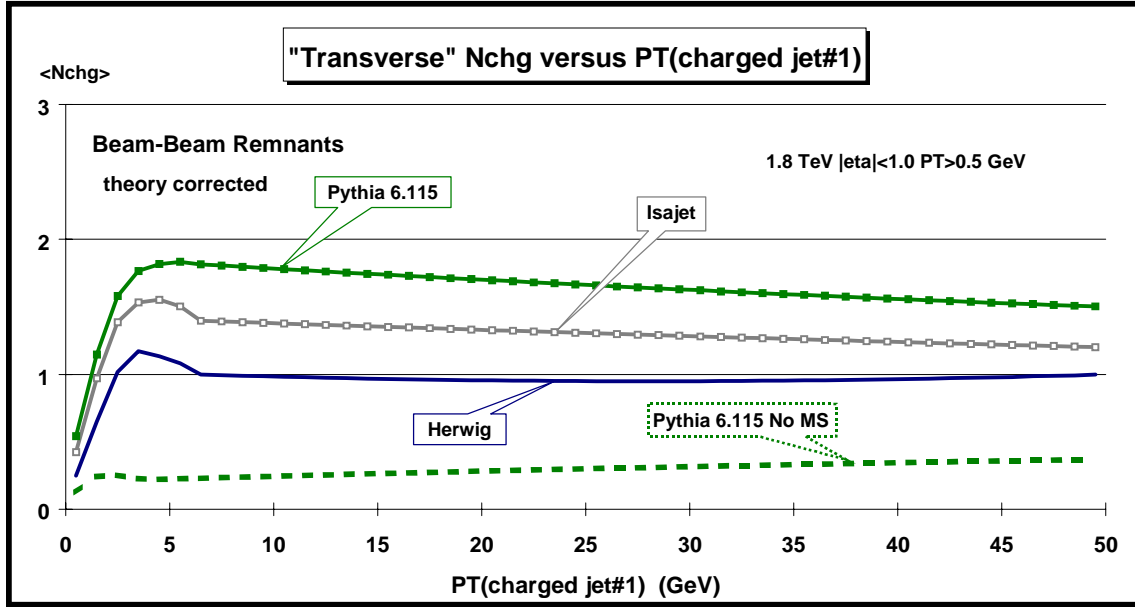
Fig. 28. Dijet data from Fig. 5 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet\#1})$  (leading charged jet) for the “transverse” region defined in Fig. 4 compared with the QCD “hard scattering” Monte-Carlo predictions of ISAJET 7.32. The predictions of ISAJET are divided into three categories: charged particles that arise from the break-up of the beam and target (beam-beam remnants), charged particles that arise from initial-state radiation, and charged particles that result from the outgoing jets plus final-state radiation (see Fig. 1). The errors on the (uncorrected) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (statistical plus systematic) of around 5%.



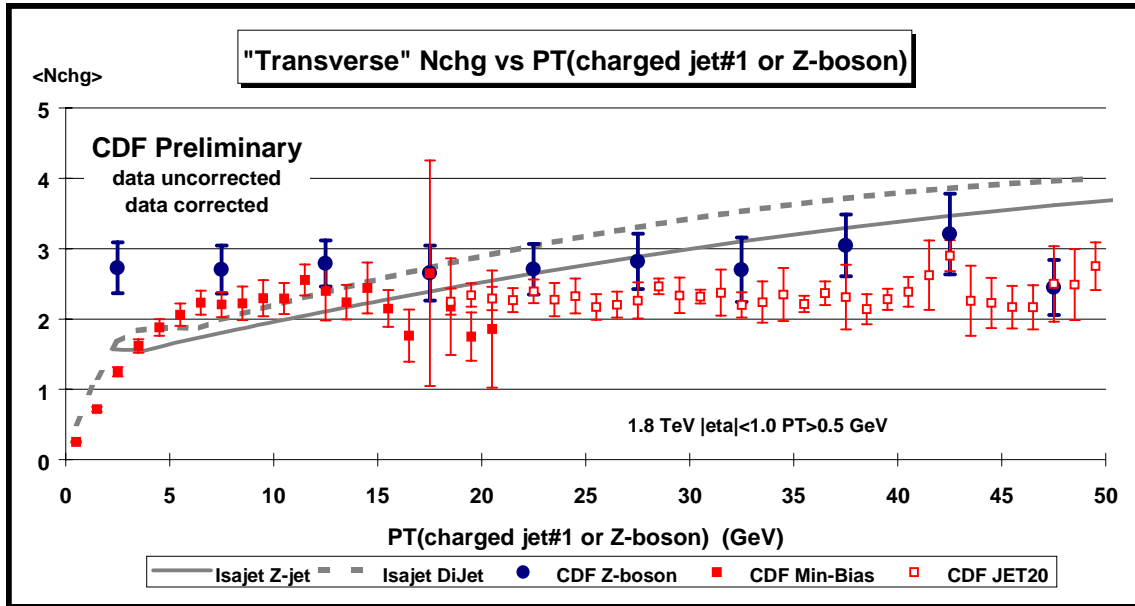
**Fig. 29.** Dijet data from Fig. 5 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet}\#1)$  (*leading charged jet*) for the “transverse” region defined in Fig. 4 compared with the QCD “hard scattering” Monte-Carlo predictions of PYTHIA 6.115. The predictions of PYTHIA are divided into two categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), and charged particles that result from the outgoing jets plus initial and final-state radiation (*hard scattering component*). For PYTHIA the beam-beam remnants include contributions from multiple parton scattering (see Fig. 2). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 30.** QCD “hard scattering” Monte-Carlo dijet predictions from HERWIG 5.9, ISAJET 7.32, and PYTHIA 6.115 of the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet}\#1)$  (*leading charged jet*) for the “transverse” region defined in Fig. 4 arising from the outgoing jets plus initial and final-state radiation (*hard scattering component*). The curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

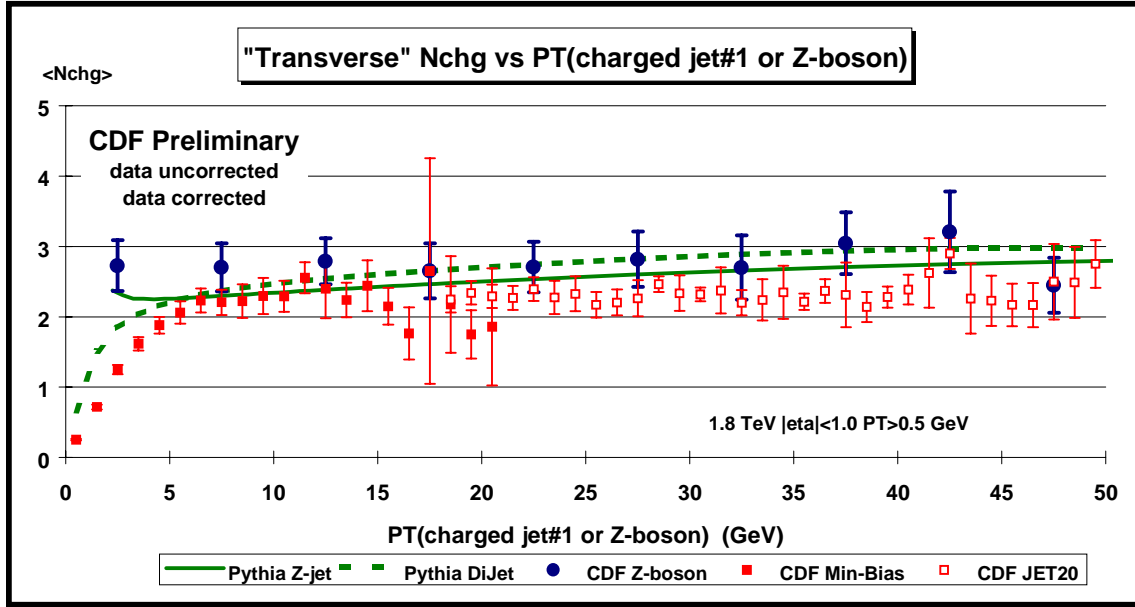


**Fig. 31.** QCD “hard scattering” Monte-Carlo dijet predictions from HERWIG 5.9, ISAJET 7.32, and PYTHIA 6.115 for the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(\text{jet\#1})$  (leading charged jet) for the “transverse” region defined in Fig. 4 arising from the break-up of the beam and target (*beam-beam remnants*). For PYTHIA the beam-beam remnants include contributions from multiple parton scattering (see Fig. 2). The curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

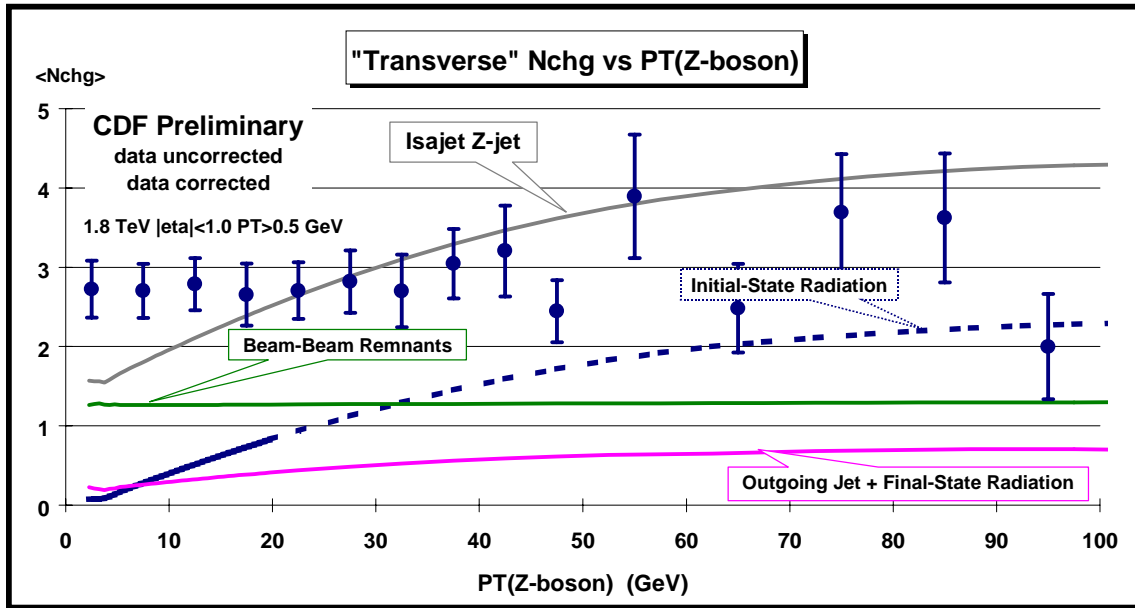


**Fig. 32.** Comparison of the dijet data from Fig. 5 and the Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “transverse” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of ISAJET 7.32 for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.

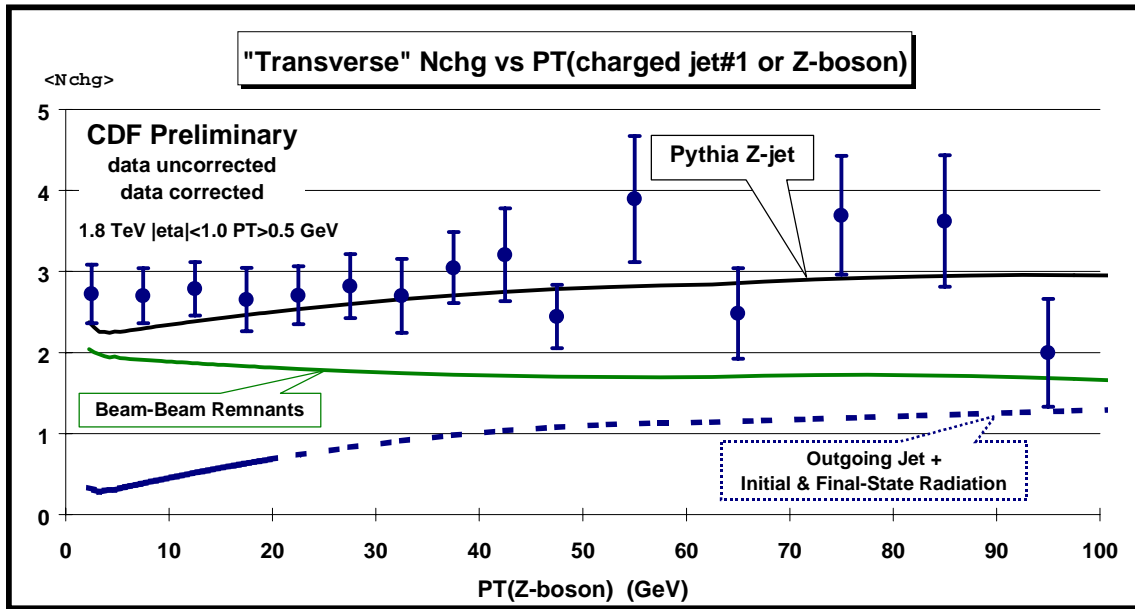




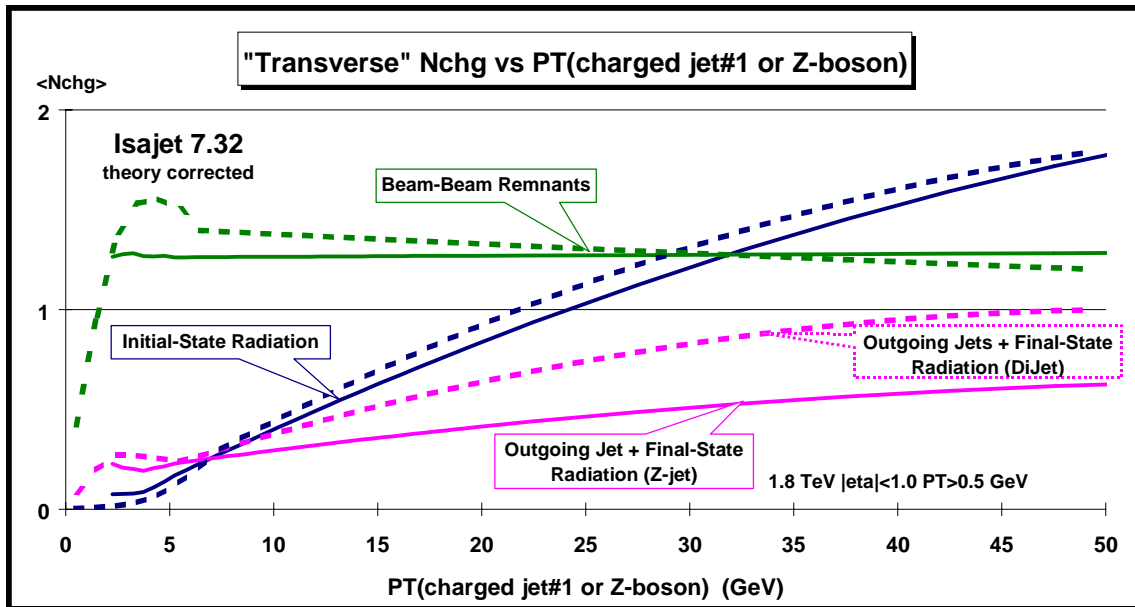
**Fig. 33.** Comparison of the dijet data from Fig. 5 and the Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “transverse” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of PYTHIA for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



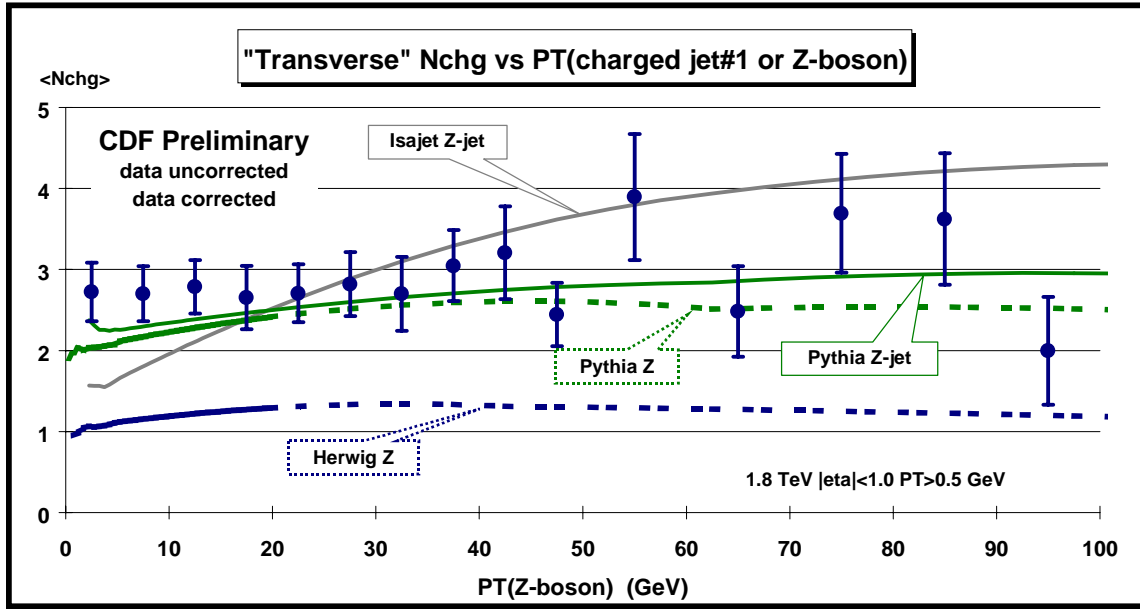
**Fig. 34.** Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $PT(Z)$  for the “transverse” region defined in Fig. 4 compared with the “Z-jet” QCD Monte-Carlo predictions of ISAJET 7.32. The predictions of ISAJET are divided into three categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), charged particles that arise from initial-state radiation, and charged particles that result from the outgoing jet plus final-state radiation (see Fig. 2). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



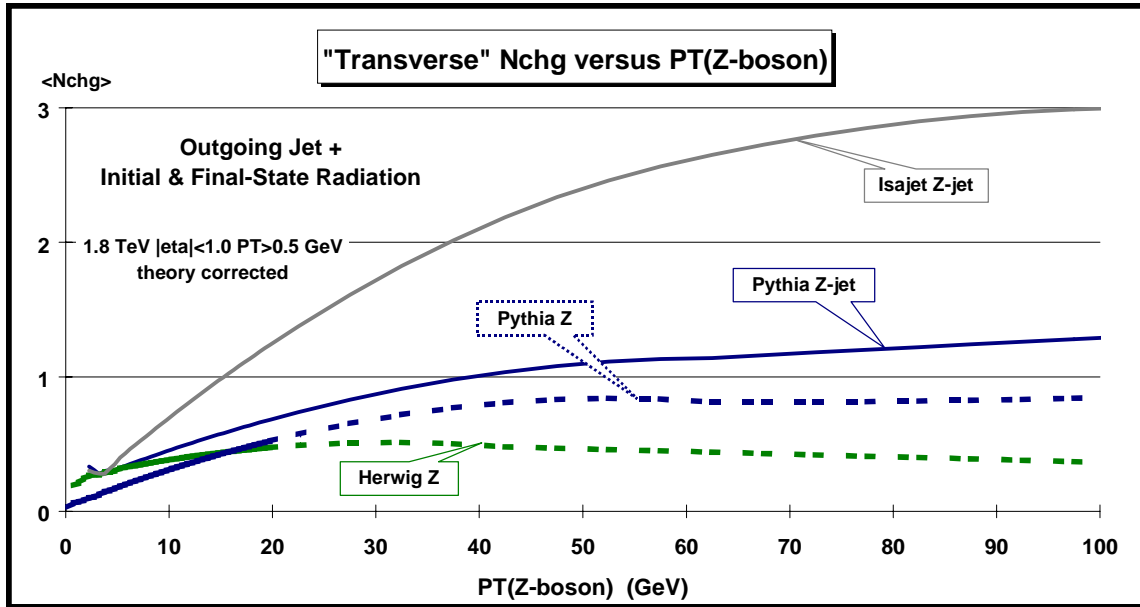
**Fig. 35.** Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “transverse” region defined in Fig. 4 compared with the “Z-jet” QCD Monte-Carlo predictions of PYTHIA 6.115. The predictions of PYTHIA are divided into two categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), and charged particles that result from the outgoing jet plus initial and final-state radiation (*hard scattering component*). For PYTHIA the beam-beam remnants include contributions from multiple parton scattering (see Fig. 2). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



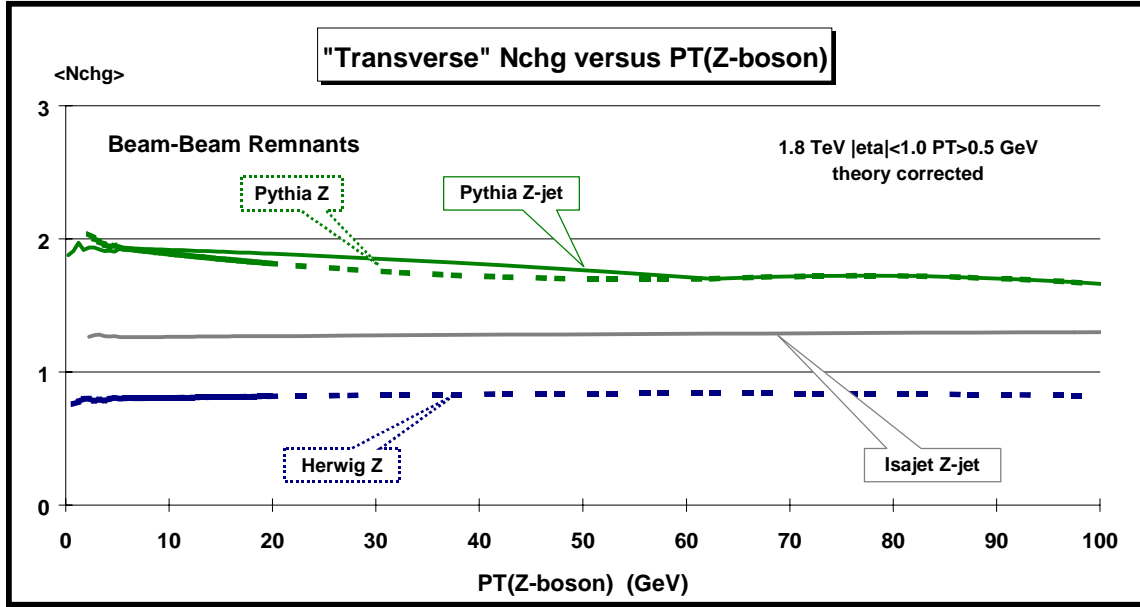
**Fig. 36.** Comparison of the QCD Monte-Carlo predictions of ISAJET 7.32 for the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “transverse” region defined in Fig. 4 for dijet (dashed) and “Z-jet” (solid) production. The predictions of ISAJET are divided into three categories: charged particles that arise from the break-up of the beam and target (*beam-beam remnants*), charged particles that arise from initial-state radiation, and charged particles that result from the outgoing jets plus final-state radiation (see Fig. 1 and Fig. 3). The curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



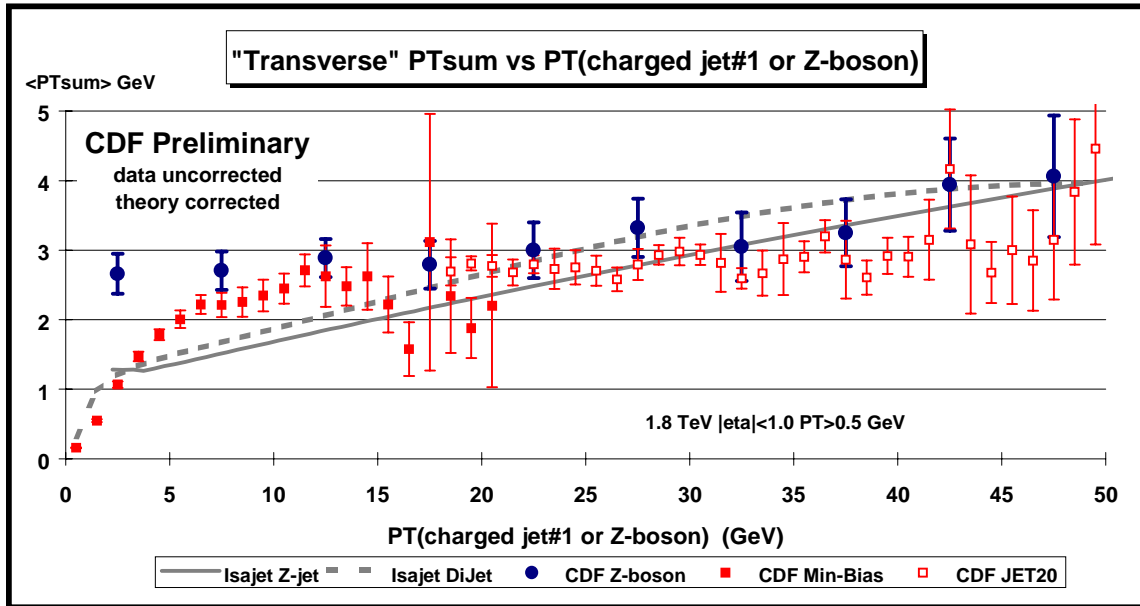
**Fig. 37.** Z-boson data from Fig. 7 on the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “transverse” region defined in Fig. 4 compared with the QCD Monte-Carlo predictions of HERWIG 5.9 (“Z”), ISAJET 7.32 (“Z-jet”), and PYTHIA 6.115 (“Z”, “Z-jet”) (see table 3). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



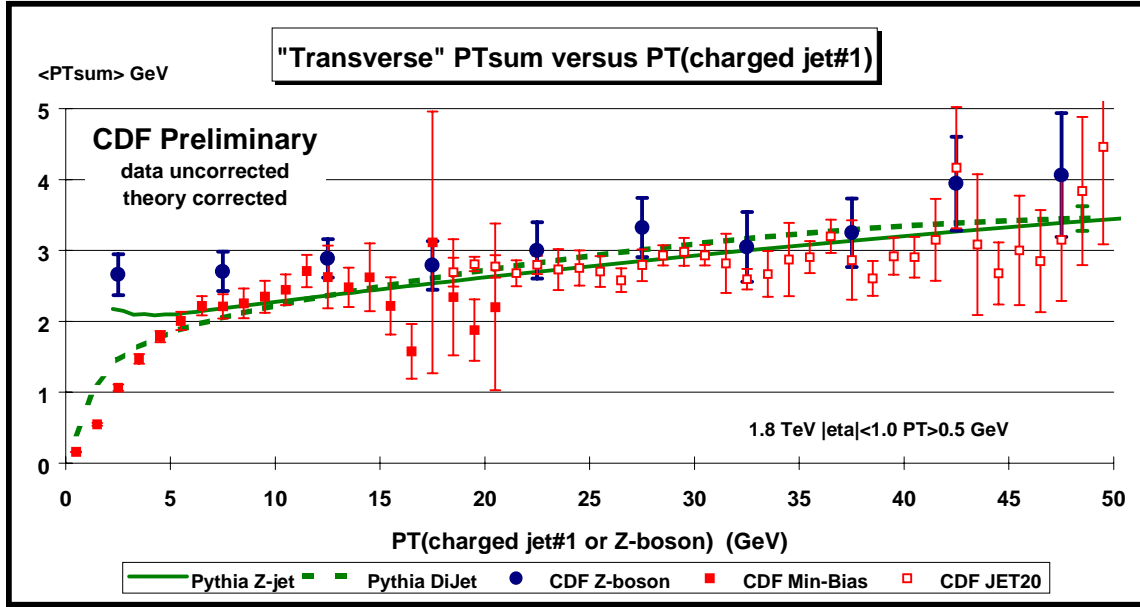
**Fig. 38.** QCD Monte-Carlo predictions from HERWIG 5.9 (“Z”), ISAJET 7.32 (“Z-jet”), and PYTHIA 6.115 (“Z”, “Z-jet”) of the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “transverse” region defined in Fig. 4 arising from the outgoing jet plus initial & final-state radiation (*hard scattering component*). The curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



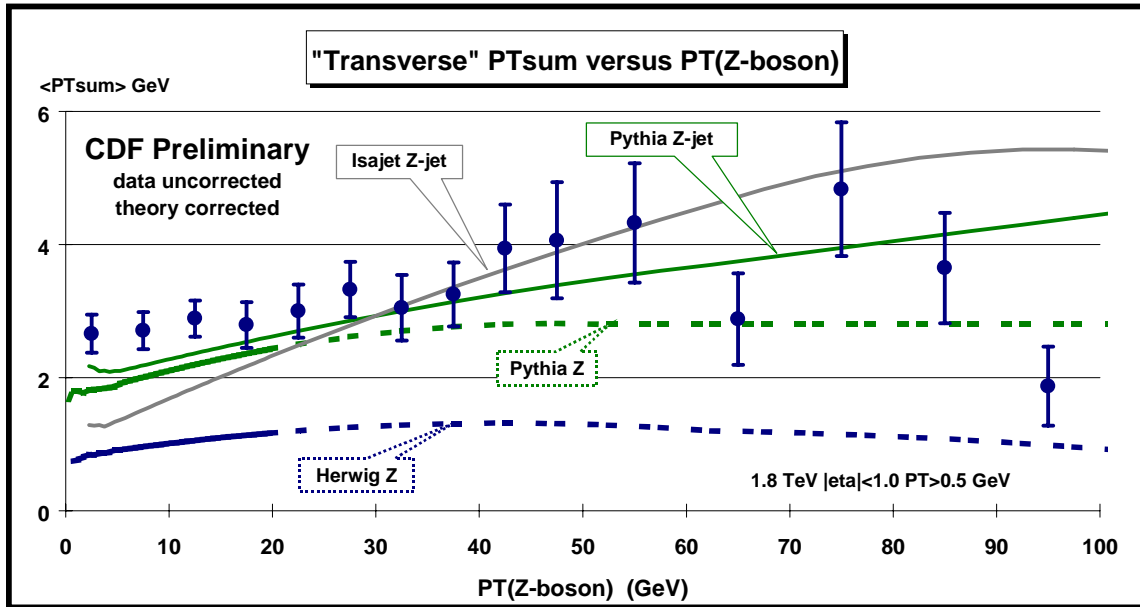
**Fig. 39.** QCD Monte-Carlo predictions from HERWIG 5.9 (“Z”), ISAJET 7.32 (“Z-jet”), and PYTHIA 6.115 (“Z”, “Z-jet”) of the average number of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “transverse” region defined in Fig. 4 arising from the break-up of the beam and target (*beam-beam remnants*). For PYTHIA the beam-beam remnants include contributions from multiple parton scattering (see Fig. 2). The curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 40.** Comparison of the dijet data from Fig. 6 and the Z-boson data from Fig. 8 on the average *scalar* PT sum of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “transverse” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of ISAJET for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 41.** Comparison of the dijet data from Fig. 6 and the Z-boson data from Fig. 8 on the average *scalar*  $P_T$  sum of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) for the “transverse” region defined in Fig. 4. The plot shows the QCD Monte-Carlo predictions of HERWIG for dijet (dashed) and “Z-jet” (solid) production. The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.



**Fig. 42.** Z-boson data from Fig. 8 on the average *scalar*  $P_T$  sum of charged particles ( $P_T > 0.5$  GeV and  $|\eta| < 1$ ) as a function of  $P_T(Z)$  for the “transverse” region defined in Fig. 4 compared with the QCD Monte-Carlo predictions of HERWIG 5.9 (“Z”), ISAJET 7.32 (“Z-jet”), and PYTHIA 6.115 (“Z”, “Z-jet”) (see table 3). The errors on the (*uncorrected*) data include both statistical and correlated systematic uncertainties. The theory curves are corrected for the track finding efficiency and have an error (*statistical plus systematic*) of around 5%.